

RAGHAV..

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$$1) \frac{d\alpha}{d\ln\mu} = -2\alpha \left(\left(\frac{11C_A - 2N_f}{3} \right) \frac{\alpha}{4\pi} + b_2 \left(\frac{\alpha}{4\pi} \right)^2 + \dots \right) \quad \text{--- (1)}$$

where $\alpha = \frac{g^2}{(4\pi)^2}$

$$\left. \begin{aligned} C_A &= 3 \\ C_F &= \frac{4}{3} \end{aligned} \right\} \text{for } SU(3)$$

Expression for $\alpha(\mu)$...

Writing (1) for the one-loop ;

$$\mu \frac{d\alpha}{d\mu} = -2\alpha \left(\left(\frac{11C_A - 2N_f}{3} \right) \frac{\alpha}{4\pi} \right)$$

$$\mu \frac{d\alpha}{d\mu} = \frac{-\alpha^2}{2\pi} \left(\left(\frac{11C_A - 2N_f}{3} \right) \right)$$

$$\Rightarrow \int_{\alpha_0}^{\alpha} \frac{-d\alpha}{\alpha^2 \chi} = \frac{1}{2\pi} \int_{\mu_0}^{\mu} \frac{d\mu}{\mu} \quad \checkmark$$

$$\frac{1}{\chi} \left| \frac{1}{\alpha} \right|_{\alpha_0}^{\alpha} = \frac{1}{2\pi} \ln \left(\frac{\mu}{\mu_0} \right)$$

$$\frac{1}{\chi} \left(\frac{1}{\alpha} - \frac{1}{\alpha_0} \right) = \frac{1}{2\pi} \ln \left(\frac{\mu}{\mu_0} \right)$$

$$\frac{1}{\alpha} = \frac{1}{\alpha_0} + \frac{\chi}{2\pi} \ln \left(\frac{\mu}{\mu_0} \right) \quad \checkmark$$

$$\cancel{\alpha} \Rightarrow \alpha_0 = \alpha + \frac{\chi \alpha d\alpha}{2\pi} \ln \left(\frac{\mu}{\mu_0} \right)$$

$$\Rightarrow \alpha(\mu) = \frac{\alpha_0}{1 + \alpha_0 \chi \ln(\mu/\mu_0)}$$

b) for $N=2$

$$C_A = 2$$

$$C_f = 3/4$$

β -function develops a zero when the signs of b_0 & b_1 are opposite. In this regime, the asymptotic freedom is lost.

$$34 C_A^2 < 2(5 C_A + 3 C_f) N_f$$

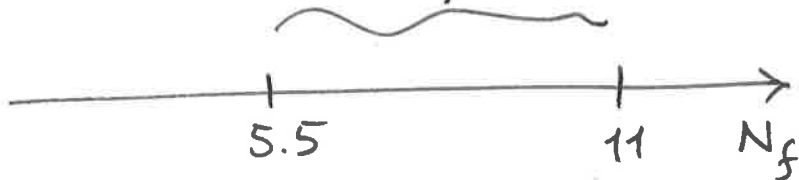
$$136 < \frac{98}{4} N_f$$

$$N_f > 5.55$$



So, when $11 > N_f > 5.55$, β -function can develop a zero away from origin.

conformal window



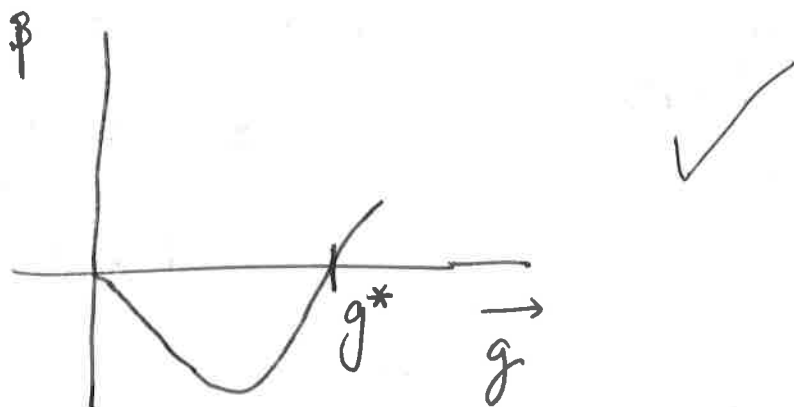
Sketch β -f^d for $N_f=0$ & $N_f > 11$

for $N_f=0$

$$b_1 = \frac{11 C_A}{3} = \frac{22}{3} ; b_2 = \left(\frac{34}{3}\right) 4 = \frac{136}{3}$$

It is interesting to note that there are no zeros of β - far away from origin in either of them since that happens for $5.5 < N_f < 11$

For ex: for $N_f = 8$



$$c) \beta(\alpha) = -\frac{b_0 \alpha^2}{4\pi} + \frac{b_1 \alpha^3}{16\pi^2}$$

$$\beta(\alpha^*) = 0$$

$$\Rightarrow \frac{b_0 \alpha^2}{4\pi} = \frac{b_1 \alpha^3}{16\pi^2 \cdot 4\pi}$$

$$\frac{b_0}{b_1} = \frac{\alpha}{4\pi}$$

$$\frac{(11C_A - 2N_f)}{34C_A^2 - 10C_A N_f - 6C_F N_f} = \frac{\alpha}{4\pi}$$

$$\frac{22 - 2N_f}{156 - 20N_f - \frac{18}{4}N_f} = \frac{\alpha}{4\pi}$$

From part (a) we see that if $\mu \rightarrow 0$

$d|_{\mu=0} \rightarrow$ some constant value.

and β -fn is zero..

This is the case of IR fixed point. The theory behaves in a quasi-conformal way.

Since, we don't have any notion of mass scale (conformal). This IR fixed point ~~is~~^{is} also known as Banks-Zaks fixed point.

~~Also,~~



$$Q^2 A_\mu^a = D_\mu^{ab} (Qc^b) - \frac{1}{2} g f^{abc} (\partial_\mu c^c c^b) - \frac{1}{2} g^2 f^{bdc} f^{cae} A_\mu^e c^d c^b \dots \textcircled{3}$$

Note that since f^{abc} is anti-symmetric, we can

recast $(\partial_\mu c^c) c^b$ as

$$(\partial_\mu c^c) c^b = \frac{1}{2} \partial_\mu (c^c c^b) \checkmark \dots \textcircled{4}$$

Using $Qc^b = -\frac{1}{2} g f^{bca} c^c c^a$ in $\textcircled{3}$ gives

$$Q^2 A_\mu^a = \frac{1}{2} g f^{abc} (\partial_\mu c^c c^b) - \frac{1}{2} g f^{abc} (\partial_\mu c^c c^b) + \frac{1}{2} g^2 f^{bdc} f^{cae} A_\mu^e c^d c^b - \frac{1}{2} g^2 f^{bdc} f^{cae} A_\mu^e c^d c^b$$

(Note: reballing of indices in the underlined term)

$= 0 \checkmark$

$$3) \gamma_A = \frac{1}{2} \frac{\partial \ln Z_3}{\partial \ln \mu} \dots \textcircled{1}$$

where Z_3 in \overline{MS} scheme is given by:

$$Z_3 = 1 + \left[\frac{5}{3} T(A) - \frac{4}{3} n_F T(R) \right] \frac{g^2}{8\pi^2} \frac{1}{\epsilon} + \mathcal{O}(g^4)$$

Rewrite $\textcircled{1}$ as

$$\gamma_A = \frac{1}{2} \frac{\partial \ln Z_3}{\partial g} \frac{\partial g}{\partial \ln \mu} \dots \textcircled{2}$$

We derived the expression for $\frac{\partial g}{\partial \ln \mu}$ to $\mathcal{O}(\epsilon)$ as

$$\frac{\partial g}{\partial \ln \mu} = -\frac{\epsilon g}{2}$$

$$\frac{\partial \ln Z_3}{\partial g} = \left[\frac{5}{3} T(A) - \frac{4}{3} n_F T(R) \right] \frac{2g}{8\pi^2 \epsilon} + \mathcal{O}(g^4)$$

$\ln(1+x) \sim x$
to $\mathcal{O}(g^2)$

Then, eq. $\textcircled{2}$ becomes

$$\begin{aligned} \gamma_A &= \frac{1}{2} \left[\frac{5}{3} T(A) - \frac{4}{3} n_F T(R) \right] \frac{2g}{8\pi^2 \epsilon} \frac{-\epsilon g}{2} \\ &= -\frac{g^2}{16\pi^2} \left[\frac{5}{3} T(A) - \frac{4}{3} n_F T(R) \right] + \mathcal{O}(g^4) \end{aligned}$$